

$i_C = I_S e^{v_{BE}/V_T} \left(1 + \frac{v_{CE}}{V_A} \right) \quad i_B = \frac{i_C}{\beta} \quad i_E = (\beta + 1) i_B$	$g_m = \frac{\partial i_C}{\partial v_{BE}} = \frac{I_C}{V_T} \text{ for a BJT}$	$r_e = \frac{V_T}{I_E} = \frac{\alpha V_T}{I_C}$	$r_\pi = \frac{V_T}{I_B} = \frac{\beta V_T}{I_C}$	$r_o = \frac{ V_A + V_{CE} }{I_C} \cong \frac{ V_A }{I_C}$
$i_D = \frac{1}{2} \mu_n \frac{W}{L} (V_{GS} - V_T)^2 \left(1 + \frac{v_{DS}}{V_A} \right) \quad V_{OV} = V_{GS} - V_T$	$g_m = \frac{2I_D}{V_{OV}} = \sqrt{2\mu_n C_{OX} \frac{W}{L} I_D} = \mu_n C_{OX} \frac{W}{L} V_{OV}$ $g_{mb} = \chi g_m \text{ for a MOSFET}$	$I_D = \frac{1}{2} k_n' \frac{W}{L} V_{OV}^2 \left(1 + \frac{v_{DS}}{V_A} \right)$ $k_n' \equiv \mu_n C_{OX} \quad \& \quad V_{OV} = V_{GS} - V_T$	$r_\pi = \infty$	$r_o = \frac{ V_A + V_{DS} }{I_D} \cong \frac{ V_A }{I_D} = \frac{1}{\lambda I_D}$ <p>where $\lambda = \frac{1}{V_A}$</p>

1) The impedance "seen" looking into the Base of a NPN/PNP can be expressed, after much algebraic pain and suffering, as shown below:

$$r_b \approx r_\pi + (\beta + 1) R_{Ext} \left(\frac{1}{1 + \left(\frac{R_{Cext}}{r_o} \right)} \right) \quad (1)$$

2) In a similar fashion, the impedance "seen" looking into the emitter of a NPN/PNP, can be expressed as shown below:

$$r_{emitter} \approx \frac{R_{Bext}}{(\beta + 1)} + r_e \left(\frac{r_o + R_{Cext}}{r_o + R_{Cext}/(\beta + 1)} \right) \quad (3)$$

3) Finally, the impedance "seen" looking into the collector of a NPN/PNP, can be expressed as, after way too much pain and suffering, as shown below:

$$r_c \approx r_o \left(1 + g_m \left(\frac{R_{Ext} \parallel (r_\pi + R_{Bext})}{1 + R_{Bext}/r_\pi} \right) \right) \quad (5)$$

4) Common Emitter (fully and partially bypassed) – the expression for the common emitter gain is represented by the following:

$$v_o/v_s = \left(\frac{v_b}{v_s} \right) \left(\frac{v_\pi}{v_b} \right) \left(\frac{v_o}{v_\pi} \right) = -\frac{R_{in}}{R_{in} + R_s} \frac{R_{Cext}}{R_{Ext} + r_e} \quad (7)$$

5) Common Base – the expression for the common base gain is represented by the following:

$$v_o/v_s = \left(\frac{v_e}{v_s} \right) \left(\frac{v_\pi}{v_e} \right) \left(\frac{v_o}{v_\pi} \right) = \frac{R_{in}}{R_{in} + R_s} g_m R_{Cext} = \frac{R_{in}}{R_{in} + R_s} \frac{R_{Cext}}{r_e} \quad (9)$$

6) Common Collector or Emitter Follower – the expression for the emitter follower is represented by the following:

$$v_o/v_s = \left(\frac{v_b}{v_s} \right) \left(\frac{v_o}{v_b} \right) = \frac{R_{in}}{R_{in} + R_s} \frac{R_{Ext}}{R_{Ext} + r_e} \quad (10)$$

Common Emitter Amplifier with R_e	Common Base Amplifier	Common Collector Amplifier
$A_v = -g_m (R_L \parallel R_c \parallel r_o) \text{ if } R_e = 0$ $A_v \cong -\frac{g_m}{1 + g_m R_e} \left[R_L \parallel R_{out} \right] \text{ if } R_e \neq 0$ $R_{in} = R_{BASE} \parallel (\beta + 1)(r_e + R_e),$ $R_{out} = R_e \parallel r_\pi + r_o \left[1 + g_m (R_e \parallel r_\pi) \right]$	$A_v = g_m \left[R_C \parallel R_L \right] \text{ if } r_o = \infty$ $R_{in} = r_e \left[\frac{R_L + r_o}{r_o + R_L/(\beta + 1)} \right]$ $R_{out} = R_e \parallel r_\pi + r_o \left[1 + g_m (R_e \parallel r_\pi) \right]$	$A_v = \frac{r_o \parallel R_L}{(r_o \parallel R_L) + r_e}$ $R_{in} = R_{BASE} \parallel (\beta + 1) \left[r_e + R_L \parallel r_o \right]$ $R_{out} = r_o \parallel R_e \parallel \left[r_e + \frac{R_{sig} \parallel R_{BASE}}{\beta + 1} \right]$
Common Source Amplifier with R_s	Common Gate Amplifier	Common Drain Amplifier
$A_v = -g_m (r_o \parallel R_L \parallel R_D) \text{ if } R_s = 0$ $A_v = -(g_m r_o) \left(\frac{R_L}{R_L + R_{out}} \right) = \frac{-g_m \left[R_L \parallel R_{out} \right]}{1 + (g_m + g_{mb}) R_s}$ $R_{in} = R_{GATE}, \quad R_{out} = r_o + \left[1 + (g_m + g_{mb}) r_o \right] R_s$	$A_v = g_m \left[R_D \parallel R_L \right] \text{ for } r_o = \infty$ $R_{in} = \frac{1}{g_m + g_{mb}} + \frac{R_D \parallel R_L}{(g_m + g_{mb}) r_o}$ $R_{out} = r_o + \left[1 + (g_m + g_{mb}) r_o \right] R_s$	$A_v = -\frac{g_m R_L'}{1 + g_m R_L'} \text{ where } R_L' = R_L \parallel r_o \parallel \left(\frac{1}{g_{mb}} \right)$ $R_{in} = R_{GATE}, \quad R_{out} = \left(\frac{1}{g_m + g_{mb}} \right) \parallel r_o \parallel R_L$

Transfer Functions and Frequency Response

1. Let
$$A(j\omega) = A_M \frac{\left(1 + j\frac{\omega}{z_1}\right)\left(1 + j\frac{\omega}{z_2}\right)\left(1 + j\frac{\omega}{z_3}\right)\cdots\left(1 + j\frac{\omega}{z_m}\right)}{\left(1 + j\frac{\omega}{p_1}\right)\left(1 + j\frac{\omega}{p_2}\right)\left(1 + j\frac{\omega}{p_3}\right)\cdots\left(1 + j\frac{\omega}{p_n}\right)}$$

a. For the Bode Magnitude Plots

$$20 \log |A(j\omega)| = 20 \log |A_M| + 20 \log \sqrt{1 + \frac{\omega^2}{z_1^2}} + 20 \log \sqrt{1 + \frac{\omega^2}{z_2^2}} + \dots + 20 \log \sqrt{1 + \frac{\omega^2}{z_m^2}} \\ - 20 \log \sqrt{1 + \frac{\omega^2}{p_1^2}} - 20 \log \sqrt{1 + \frac{\omega^2}{p_2^2}} - \dots - 20 \log \sqrt{1 + \frac{\omega^2}{p_n^2}}$$

b. For the Bode Phase Plots

$$\Phi \angle A(j\omega) = \tan^{-1} A_M + \tan^{-1} \left(\frac{\omega}{z_1} \right) + \tan^{-1} \left(\frac{\omega}{z_2} \right) + \dots + \tan^{-1} \left(\frac{\omega}{z_m} \right) \\ - \tan^{-1} \left(\frac{\omega}{p_1} \right) - \tan^{-1} \left(\frac{\omega}{p_2} \right) - \dots - \tan^{-1} \left(\frac{\omega}{p_n} \right)$$

2. Determining the 3-dB frequencies of amplifiers

(A) Poles and Zeros are Known or Can Be Easily Determined

Low-Frequency Response	High-Frequency Response
$A(s) \cong A_M F_L(s)$ $F_L(s) = \frac{\left(1 + s/\omega_{z_1}\right)\left(1 + s/\omega_{z_2}\right)\cdots\left(1 + s/\omega_{z_{nL}}\right)}{\left(1 + s/\omega_{p_1}\right)\left(1 + s/\omega_{p_2}\right)\cdots\left(1 + s/\omega_{p_{nL}}\right)}$	$A(s) \cong A_M F_H(s)$ $F_H(s) = \frac{\left(1 + s/\omega_{z_1}\right)\left(1 + s/\omega_{z_2}\right)\cdots\left(1 + s/\omega_{z_{nH}}\right)}{\left(1 + s/\omega_{p_1}\right)\left(1 + s/\omega_{p_2}\right)\cdots\left(1 + s/\omega_{p_{nH}}\right)}$
<p>If $\omega_{p_1} \gg \omega_{p_2}, \omega_{p_3}, \dots, \omega_{z_1}, \omega_{z_2}, \dots$ then for frequencies near the midband:</p> $F_L(s) \cong \frac{s}{s + \omega_{p_1}}, \text{ and } \omega_L \cong \omega_{p_1}$	<p>If $\omega_{p_1} \ll \omega_{p_2}, \omega_{p_3}, \dots, \omega_{z_1}, \omega_{z_2}, \dots$ then for frequencies near the midband:</p> $F_H(s) \cong \frac{1}{1 + s/\omega_{p_1}}, \text{ and } \omega_H \cong \omega_{p_1}$
<p>Otherwise,</p> $\omega_L \cong \sqrt{\omega_{p_1}^2 + \omega_{p_2}^2 + \dots - 2\left(\omega_{z_1}^2 + \omega_{z_2}^2 + \dots\right)}$	<p>Otherwise,</p> $\omega_H \cong \frac{1}{\sqrt{\frac{1}{\omega_{p_1}^2} + \frac{1}{\omega_{p_2}^2} + \dots - 2\left(\frac{1}{\omega_{z_1}^2} + \frac{1}{\omega_{z_2}^2} + \dots\right)}}$

(B) Poles and Zeros Cannot Be Easily Determined

Low-Frequency Response	High-Frequency Response
$F_L(s) = \frac{s^{nL} + d_1 s^{nL-1} + \dots}{s^{nL} + e_1 s^{nL-1} + \dots}$	$F_H(s) = \frac{1 + a_1 s + a_2 s^2 + \dots}{1 + b_1 s + b_2 s^2 + \dots}$
$e_1 = \omega_{p_1} + \omega_{p_2} + \dots + \omega_{p_{nL}} = \sum_{i=1}^{nL} \frac{1}{C_i R_{i_s}}$	$b_1 = \frac{1}{\omega_{p_1}} + \frac{1}{\omega_{p_2}} + \dots + \frac{1}{\omega_{p_{nH}}} = \sum_{i=1}^{nH} C_i R_{i_o}$
<p>If a dominant pole exists (say, P_1), then</p> $\omega_L \cong \sum_{i=1}^{nL} \frac{1}{C_i R_{i_s}}$	<p>If a dominant pole exists (say, P_1), then</p> $\omega_H \cong 1 / \sum_{i=1}^{nH} C_i R_{i_o}$